From the Schrödinger problem to the Monge-Kantororich problem

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Outline

- Monge, Kantorovich and Schrödinger
 - Optimal transport
 - Schrödinger problem
- 2 From entropy to transport
 - Freezing the particle system
 - Optimal transport is the frozen limit
- 3 Some geometry

The Monge transport problem

Take

- X, a measurable space;
- $c: \mathcal{X} \times \mathcal{X} \to [0, \infty)$, a cost function;
- $\mu_0, \mu_1 \in \mathcal{P}(\mathcal{X})$, two probability measures on \mathcal{X} .

Consider the measurable maps $T: x \in \mathcal{X} \mapsto y = T(x) \in \mathcal{X}$ such that

$$\mu_1 = T_{\#}\mu_0$$
 (image measure).

Monge problem

minimize
$$T \mapsto \int_{\mathcal{X}} c(x, T(x)) \, \mu_0(dx)$$
 subject to $\mu_1 = T_{\#}\mu_0$. (M)

Example:
$$\mathcal{X} = \mathbb{R}^d$$
, $c(x, y) = ||y - x|| \text{ or } ||y - x||^2$.

The Monge-Kantorovich transport problem

We want to transport μ_0 onto μ_1 with a minimal cost.

A transport plan is
$$\pi \in \mathcal{P}(\mathcal{X} \times \mathcal{X})$$
 such that $\left\{ \begin{array}{l} \text{first marginal} & := \pi_0 = \mu_0 \\ \text{second marginal} & := \pi_1 = \mu_1 \end{array} \right.$
The joint law $\pi = \mathcal{L}(X_0, X_1)$ is a *coupling* of $\mu_0 = \mathcal{L}(X_0)$ and $\mu_1 = \mathcal{L}(X_1)$. With $X_0, X_1 : \mathcal{X} \times \mathcal{X} \to \mathcal{X}$ the first and second projections:

$$X_{0\#}\pi = \mu_0, \qquad X_{1\#}\pi = \mu_1.$$

Monge-Kantorovich problem

minimize
$$\pi \in \mathcal{P}(\mathcal{X} \times \mathcal{X}) \mapsto \int_{\mathcal{X} \times \mathcal{X}} c(x, y) \, \pi(dxdy)$$
 subject to $X_{0\#}\pi = \mu_0, X_{1\#}\pi = \mu_1.$ (MK)

It is a relaxed version of the Monge problem.

Disintegration:
$$\pi(dxdy) = \pi_0(dx)\pi(dy|X_0 = x)$$
.
Monge corresponds to $\pi(dy|X_0 = x) = \delta_{T(x)}(dy)$.

Wasserstein metric

Take (\mathcal{X}, d) a metric space.

Consider the cost $c(x, y) = d^{p}(x, y), p = 1 \text{ or } 2.$

The Wasserstein metric on $\mathcal{P}_{p}(\mathcal{X})$ is

$$\mathit{W}_{p}(\mu_{0},\mu_{1}) = \inf \left\{ \left(\mathrm{MK} \right)_{\mathit{dP}} \right\}^{1/p}, \quad \mu_{0},\mu_{1} \in \mathcal{P}(\mathcal{X}).$$

Result: W_{ρ} is a metric on $\mathcal{P}_{\rho}(\mathcal{X}) := \{ \mu \in \mathcal{P}(\mathcal{X}); \int_{\mathcal{X}} d^{\rho}(x_o, x) \, \mu(dx) < \infty \}.$

The Schrödinger problem

Take *n* independent Brownian particles in $\mathcal{X} = \mathbb{R}^d$:

$$X_i(t) = x_i + \sigma W_i(t), \quad 0 \le t \le 1, i = 1, \ldots, n.$$

- σ^2 is the temperature of the heat bath;
- W_1, \ldots, W_n are independent Brownian motions;
- \bullet x_1, \ldots, x_n are the initial positions.

Suppose that: $L_0^n := \frac{1}{n} \sum_{i=1}^n \delta_{X_i(0)} = \frac{1}{n} \sum_{i=1}^n \delta_{X_i} \underset{n \to \infty}{\longrightarrow} \mu_0 \text{ in } \mathcal{P}(\mathcal{X}).$ Then:

$$L^n := \frac{1}{n} \sum_{i=1}^n \delta_{X_i} \underset{n \to \infty}{\longrightarrow} R \quad \text{in} \quad \mathcal{P}(\Omega)$$

where $\Omega = \{\text{paths}\} = C([0, 1], \mathcal{X}) \text{ and } R := \mathcal{L}(X_0 + \sigma W), \ W \perp X_0 \sim \mu_0.$

The Schrödinger problem

In particular, with $R_1 = \mu_0 * \mathcal{N}(0, Id)$,

$$L_1^n := \frac{1}{n} \sum_{i=1}^n \delta_{X_i(1)} \underset{n \to \infty}{\to} R_1 := X_{1\#}R \text{ in } \mathcal{P}(\mathcal{X})$$

With $n = \infty$ the system starts from $R_0 = \mu_0$ and ends up at R_1 almost surely. But n is finite.

Schrödinger's question

Suppose that at time t=1 you observe L_1^n near μ_1 , far from R_1 . What is the most likely path $t\in[0,1]\mapsto L_t^n\in\mathcal{P}(\mathcal{X})$ of your particle system?

This is a large deviation problem.

Schrödinger's answer

Relative entropy: $H(P|R) := \int_{\Omega} \log(\frac{dP}{dR}) dP \in [0, \infty], P, R \in \mathcal{P}(\Omega).$

Theorem (Sanov's theorem)

For
$$A \in \mathcal{P}(\Omega)$$
, $\mathbb{P}(L^n \in A) \underset{n \to \infty}{\simeq} \exp \{-n \inf_{P \in A} \frac{H(P|R)}{H(P|R)}\}$.

Conditionally on $L_0^n \underset{n \to \infty}{\to} \mu_0$ and $L_1^n \underset{n \to \infty}{\to} \mu_1$,

$$L^n \underset{n \to \infty}{\longrightarrow} \widehat{P}$$
 almost surely in $\mathcal{P}(\Omega)$

where \widehat{P} is the unique solution of

Schrödinger dynamical problem

minimize
$$P \in \mathcal{P}(\Omega) \mapsto \mathcal{H}(P|R)$$

subject to $X_{0\#}P = \mu_0, X_{1\#}P = \mu_1.$ (S)

Schrödinger's answer

Denote for all $P \in \mathcal{P}(\Omega)$,

- $P_t = X_{t\#}P \in \mathcal{P}(\mathcal{X})$ the law of position at time t;
- $P_{01} = (X_0, X_1)_{\#} P \in \mathcal{P}(\mathcal{X} \times \mathcal{X})$ the joint law of initial and final positions;
- $P^{xy} = P(\cdot | X_0 = x, X_1 = y)$ the bridge between x and y;
- $\bullet R^{\pi}(\cdot) := \int_{\mathcal{X} \times \mathcal{X}} R^{xy}(\cdot) \pi(dxdy).$

Tensorization: $H(P|R) = H(P_{01}|R_{01}) + \int_{X \times X} H(P^{xy}|R^{xy}) P_{01}(dxdy)$.

When $P_{01} = \pi \in \mathcal{P}(\mathcal{X} \times \mathcal{X})$ is fixed, this is minimal when $P^{xy} = R^{xy}$ for all $x, y \in \mathcal{X}$. That is for $P = \mathbb{R}^{\pi}$.

The solution of (S) is

$$\hat{P} = R^{\hat{\pi}}$$

where $\hat{\pi}$ is the solution of

Schrödinger problem

minimize
$$\pi \in \mathcal{P}(\mathcal{X} \times \mathcal{X}) \mapsto H(\pi|R_{01})$$

subject to $X_{0\#}\pi = \mu_0, X_{1\#}\pi = \mu_1.$ (S_{01})

Schilder's theorem

The main idea is to let the temperature $\epsilon = \sigma^2$ tend down to zero.

Let
$$X^{\epsilon}(t) = X^{\epsilon}(0) + \sqrt{\epsilon}W_t$$
, $0 \le t \le 1$.

For all path
$$\omega = (\omega_t)_{0 \le t \le 1}$$
, denote $C(\omega) = \int_0^1 \frac{1}{2} |\dot{\omega}_t|^2 dt \in [0, \infty]$.

Theorem (Schilder)

For all
$$A \subset \Omega$$
, $\mathbb{P}(X^{\epsilon} \in A) \underset{\epsilon \to 0}{\asymp} \exp\left\{-\frac{1}{\epsilon}\inf_{\omega \in A}[I_0(\omega_0) + C(\omega)]\right\}$.

Conditionally on $X_0^{\epsilon} \xrightarrow{} x$ and $X_1^{\epsilon} \xrightarrow{} y$,

 X^{ϵ} tends almost surely to the solution of the

Geodesic problem

minimize
$$\omega \in \Omega \mapsto C(\omega) = \int_0^1 \frac{1}{2} |\dot{\omega}_t|^2 dt$$
 subject to $\omega_0 = x, \omega_1 = y$.

Solution: γ^{xy} , the constant speed straight line between x and y and

$$c(x, y) := \inf\{C(\omega); \omega_0 = x, \omega_1 = y\} = C(\gamma^{xy}) = ||y - x||^2/2$$

is the popular quadratic cost

Schrödinger problem in a cold world

We mix Sanov and Schilder's theorems.

- $R^{\epsilon} \in \mathcal{P}(\Omega)$ is the law of X^{ϵ} ;
- $X_1^{\epsilon}, \dots, X_n^{\epsilon}$ are independent copies of X^{ϵ} ;
- $L^{n,\epsilon} = \frac{1}{n} \sum_{i=1}^{n} \delta_{X_{i}^{\epsilon}}$ is a random element in $\mathcal{P}(\Omega)$.

Theorem (Dawson-Gärtner, L.)

Assume that $\{X^{\epsilon}; \epsilon>0\}$ obeys the LDP: $R^{\epsilon}(\cdot)\underset{\epsilon\to 0}{\asymp} \exp\left\{-\frac{1}{\epsilon}C(\cdot)\right\}$. Then,

$$\Gamma$$
- $\lim_{\epsilon \to 0} \epsilon H(P|R^{\epsilon}) = \int_{\Omega} C dP \in [0, \infty]$

and $\{L^{n,\epsilon}; \epsilon > 0, n \ge 1\}$ obeys the double index LDP

$$\mathbb{P}(L^{n,\epsilon} \in A) \underset{\epsilon \to 0, n \to \infty}{\asymp} \exp\left\{-\frac{n}{\epsilon} \inf_{P \in A} \int_{\Omega} C \, dP\right\}.$$

Idea of the proof

$$\epsilon H(P|R^{\epsilon}) = \epsilon \sup_{f} \left\{ \int f \, dP - \log \int e^{f} \, dR^{\epsilon} \right\}$$

$$= \sup_{f} \left\{ \epsilon \int f / \epsilon \, dP - \epsilon \log \int e^{f/\epsilon} \, dR^{\epsilon} \right\}$$

$$= \sup_{f} \left\{ \int f \, dP - \epsilon \log \int e^{f/\epsilon} \, dR^{\epsilon} \right\}$$
But $R^{\epsilon}(\cdot) \underset{\epsilon \to 0}{\approx} \exp\{-C(\cdot)/\epsilon\}$ gives $\int e^{f/\epsilon} \, dR^{\epsilon} \underset{\epsilon \to 0}{\approx} \sup_{\Omega} \exp\{(f-C)/\epsilon\}$ and
$$\epsilon \log \int e^{f/\epsilon} \, dR^{\epsilon} \underset{\epsilon \to 0}{\to} \sup_{\Omega} (f-C).$$

$$\epsilon H(P|R^{\epsilon}) \underset{\epsilon \to 0}{\Longrightarrow} \sup_{f} \left\{ \int f \, dP - \sup_{f} (f-C) \, dP \right\}$$

$$= \sup_{f} \left\{ \int C \, dP + \int [(f-C) - \sup_{f} (f-C)] \, dP \right\}$$

$$= \int_{C} C \, dP.$$
Cy. Légand Entropy and Optimal Transport

Γ-convergence

Lower semicontinuous envelope: $\operatorname{lsc}(h)(x) := \sup_{V \in \mathcal{V}(x)} \inf_{y \in V} h(y)$. Γ -limit: Γ - $\lim_{\epsilon \to 0} h^{\epsilon}(x) := \sup_{V \in \mathcal{V}(x)} \lim_{\epsilon \to 0} \inf_{y \in V} h^{\epsilon}(y)$. This implies that there exists a sequence $\{x^{\epsilon}\}$ such that

exists a sequence (x) such the

$$x^{\epsilon} \underset{\epsilon \to 0}{\longrightarrow} x; \quad h^{\epsilon}(x^{\epsilon}) \underset{\epsilon \to 0}{\longrightarrow} h(x).$$

More.

$$\min h^{\epsilon} \underset{\epsilon \to 0}{\longrightarrow} \min h; \quad \operatorname{argmin} h^{\epsilon} \underset{\epsilon \to 0}{\longrightarrow} \operatorname{argmin} h.$$

In particular, Γ - $\lim_{\epsilon\to 0} \epsilon H(P|R^\epsilon) = \int_{\Omega} C \, dP$ implies that there exists a sequence $\{P^\epsilon\}$ such that

$$P^{\epsilon} \underset{\epsilon \to 0}{\longrightarrow} P; \quad \epsilon H(P^{\epsilon}|R^{\epsilon}) \underset{\epsilon \to 0}{\longrightarrow} \int_{\Omega} C dP.$$

Remark: $\{R^{\epsilon}\}$ is a family of mutually singular measures.

Contraction principle

I skip the technical details about the initial condition.

Idea: Take R_0 as the reversing measure of the Markov dynamics.

Here R_0 is Lebesgue measure, which is unbounded.

By the contraction principle

$$R^{\epsilon}(d\omega) \underset{\epsilon \to 0}{symp} \exp \left\{ -\frac{1}{\epsilon} C(\omega) \right\}$$

gives

$$R_{01}^{\epsilon}(dxdy) \underset{\epsilon \to 0}{\approx} \exp\left\{-\frac{1}{\epsilon}c(x,y)\right\}$$

with

$$\begin{split} c(x,y) &= \inf\{C(\omega); \omega : \omega_0 = x, \omega_1 = y\} \\ &= \inf\left\{\int_0^1 \frac{1}{2} \|\dot{\omega}_t\|^2 \, dt; \omega : \omega_0 = x, \omega_1 = y\right\} \end{split}$$

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Γ-convergence is stable under contraction.

Limit of the minimization problems

Theorem (Dynamical problems)

In the space $\mathcal{P}(\Omega)$, we have the following convergence

$$\begin{array}{lll} (\mathcal{S}^{\epsilon}): & \textit{minimize } P \mapsto \epsilon H(P|R^{\epsilon}) & \textit{subject to} & P_0 = \mu_0^{\epsilon}, & P_1 = \mu_1^{\epsilon} \\ \Gamma\text{-}\lim_{\epsilon} \downarrow & & \downarrow & \downarrow \\ (T): & \textit{minimize } P \mapsto \int_{\Omega} \textit{C dP} & \textit{subject to} & P_0 = \mu_0, & P_1 = \mu_1 \end{array}$$

And by contraction, we obtain the

Theorem (From (S) to (MK))

In the space $\mathcal{P}(\mathcal{X} \times \mathcal{X})$, we have the following convergence

$$\begin{array}{cccc} (S_{01}^{\epsilon}): & \textit{minimize } \pi \mapsto \epsilon \textit{H}(\pi | R_{01}^{\epsilon}) & \textit{subject to} & \pi_0 = \mu_0^{\epsilon}, & \pi_1 = \mu_1^{\epsilon} \\ \Gamma\text{-lim}_{\epsilon} \downarrow & & \downarrow & \downarrow \\ (\textit{MK}): & \textit{minimize } \pi \mapsto \int_{\mathcal{X} \times \mathcal{X}} \textit{C} \, \textit{d} \pi & \textit{subject to} & \pi_0 = \mu_0, & \pi_1 = \mu_1^{\epsilon} \\ \end{array}$$

where

$$R_{01}^{\epsilon}(dxdy)\underset{\epsilon \to 0}{symp} \exp\left\{-rac{1}{\epsilon}c(x,y)
ight\}.$$

Limit of the minimizers

Theorem (Convergence of the minimizers)

There exists a sequence $(\mu_0^{\epsilon}, \mu_1^{\epsilon}) \underset{k \to \infty}{\to} (\mu_0, \mu_1)$ such that

$$\begin{array}{cccc} \widehat{P}^{\epsilon}(\cdot) & = & \int_{\mathcal{X}\times\mathcal{X}} R^{\epsilon,xy}(\cdot)\, \hat{\pi}^{\epsilon}(\textit{dxdy}); & \widehat{P}^{\epsilon}_{0} = \hat{\pi}^{\epsilon}_{0} = \mu^{\epsilon}_{0}, \widehat{P}^{\epsilon}_{1} = \hat{\pi}^{\epsilon}_{1} = \mu^{\epsilon}_{1} \\ \downarrow & & \downarrow \\ \widehat{P}(\cdot) & = & \int_{\mathcal{X}\times\mathcal{X}} \delta_{\gamma^{xy}}(\cdot)\, \hat{\pi}(\textit{dxdy}); & \widehat{P}_{0} = \hat{\pi}_{0} = \mu_{0}, \widehat{P}_{1} = \hat{\pi}_{1} = \mu_{1} \end{array}$$

Recall: $R^{\epsilon,xy}$ is the bridge of R^{ϵ} and γ^{xy} is the geodesic between x and y.

This convergence is a tool for approximating some features of geometry by means of probability theory.

Other dynamics

For any $k \geq 1$, let R^k be the law of the random walk in $\mathcal{X} = \mathbb{R}^d$

$$X^{k}(t) = X^{k}(0) + \frac{1}{k} \sum_{j=1}^{\lfloor kt \rfloor} Z_{j}, \quad 0 \leq t \leq 1$$

where Z_1, \ldots, Z_k are independent copies of Z. Define

- $C(\omega) = \int_0^1 c(\dot{\omega}_t) dt \in [0, \infty], \quad \omega \in \Omega,$
- $c(v) = \sup_{p \in \mathbb{R}^d} \{ v \cdot p \log \mathbb{E} e^{p \cdot Z} \}, \quad v \in \mathbb{R}^d.$

Theorem

- ① Γ- $\lim_{k\to\infty} \frac{1}{k} H(P|R^k) = \int_{\Omega} C dP$, $P \in \mathcal{P}(\Omega)$,
- 2 Γ $\lim_{k\to\infty} \frac{1}{k} H(\pi|R_{01}^k) = \int_{\Omega} c(y-x) \pi(dxdy), \quad \pi \in \mathcal{P}(\mathcal{X} \times \mathcal{X})$

A geometry on $\mathcal{P}(\mathcal{X})$

McCann, Gangbo, Otto, Villani,...

Definition ("Geodesics" in $\mathcal{P}(\mathcal{X})$)

The flow

$$t\mapsto \mu_t:=\widehat{P}_t(\cdot)=\int_{\mathcal{X} imes\mathcal{X}}\delta_{\gamma^{xy}(t)}(\cdot)\,\widehat{\pi}(extit{dxdy})\in\mathcal{P}(\mathcal{X})$$

where $\hat{\pi}$ is a solution of (MK) is a kind of geodesic from μ_0 to μ_1 in $\mathcal{P}(\mathcal{X})$. We write it

$$\mu_t = [\mu_0, \mu_1]_t, \quad 0 \le t \le 1$$

This "geodesic" is called a *displacement interpolation*. It is approximated by the entropy minimizers

$$t\mapsto \mu^\epsilon_t:=\widehat{P}^\epsilon_t(\cdot)=\int_{\mathcal{X} imes\mathcal{X}} R^{\epsilon, extit{xy}}_t(\cdot)\,\hat{\pi}^\epsilon(extit{dxdy})\in \mathcal{P}(\mathcal{X}).$$

Remark: (MK) might admit infinitely many solutions. (S_{01}^{ϵ}) admits a unique solution.

We select the viscosity solution.

Ricci curvature

Active field of research: Lott, Sturm, von Renesse, Villani,...

Theorem

Let $\mathcal X$ be a connected Riemannian manifold. Then, $\mathcal X$ has a nonnegative Ricci curvature if and only if the relative entropy $H(\cdot|\mathrm{vol})$ is displacement convex.

That is, for any displacement interpolation $(\mu_t)_{0 \le t \le 1}$,

$$H(\mu_t|\text{vol}) \leq (1-t)H(\mu_0|\text{vol}) + tH(\mu_1|\text{vol}), \forall 0 \leq t \leq 1.$$

This theory works mainly with the quadratic transport.

This gives a very interesting definition of *lower-bounded-Ricci-curvature* in a general metric space \mathcal{X} .

Work in progress

Try to obtain curvature properties of the measure space (\mathcal{X}, ρ) by means of the second derivative of the function

$$t\mapsto H(\widehat{P}_t|\rho)$$

where

- with respect to a general Markov reversible Markov process R with reversing measure ρ.

Take advantage of the convergence of the functions

$$t\mapsto H(\widehat{P}^{\epsilon}_t|\rho^{\epsilon})$$

where

- \widehat{P}^{ϵ} is the Schrödinger minimizer
- with respect to a general Markov reversible Markov process R^{ϵ} with reversing measure $\rho^{\epsilon} \xrightarrow[\epsilon \to 0]{} \rho$.

Monge, Kantorovich and Schrödinger From entropy to transport Some geometry

Thank you for your attention.