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Abstract

We investigate the differential geometry of bipartite quantum states. In particular the manifold structures of pure bipartite states are studied in detail. The manifolds with respect to all normalized pure states of arbitrarily given Schmidt ranks or Schmidt coefficients are explicitly presented. The dimensions of the related manifolds are calculated.

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Quantum entanglement constitutes the most important resource in quantum information processing such as quantum teleportation, dense coding, quantum cryptography, quantum error correction and quantum repeater [?]. The marvelous properties of quantum entanglement are from the special structures of the multipartite quantum states. Great efforts have been focused on the proper description and quantification of quantum entanglement [?], the separability [?], the equivalence of quantum states under local unitary transformations or under stochastic local operations and classical communication (SLOCC) for multipartite quantum systems [?].

The geometry of quantum states on a single vector space has been discussed in [? ?] recently. Let \mathcal{H} be an n-dimensional complex Hilbert space. The space of density matrices on \mathcal{H} , $D(\mathcal{H})$, is naturally a manifold stratified space with the stratification induced by the rank of the state. The space of all density matrices with rank r, $D^r(\mathcal{H})$, $r = 1, 2, \dots, n$, is a smooth and connected manifold of real dimension $2nr - r^2 - 1$. In particular, $D^1(\mathcal{H})$ is the set of pure states. Every element of $D(\mathcal{H})$ is a convex combination of points from $D^1(\mathcal{H})$. It is shown that $D^1(\mathcal{H})$ is a complex manifold which is isomorphic to the n-1 dimensional complex projective space, $D^1(\mathcal{H}) \simeq CP^{n-1}$, with a metric g determined by the inner product $\langle M, N \rangle = \frac{1}{2}TrMN$ for density matrices M and N. One can define the Hermitian structure h on $D^1(\mathcal{H})$ by g. In fact, by straightforward calculation, we have

$$h^{(\alpha)} = \sum_{k,j} h_{kj}^{(\alpha)} dz_k \otimes d\overline{z_j}, \quad h^{(\alpha)} = h_{|D_{\alpha}}, \quad \alpha = 1, ..., n,$$

where

$$h_{kj}^{(\alpha)} = \frac{(1 + \sum_{l=1, l \neq \alpha}^{n} |z_l|^2) \delta_{kj} - z_j \overline{z_k}}{(1 + \sum_{l=1, l \neq \alpha}^{n} |z_l|^2)},$$

 D_{α} is the α -th coordinate chart with local complex coordinates z and \overline{z} . Hence it is clear that h differs from the Fubini-Study metric on $\mathbb{C}P^{n-1}$ by a constant multiple.

The quantum entanglement concerns composite systems. In [?] the entanglement has been discussed in the view of geometry. In this paper we investigate the manifold structures and classification of pure bipartite states. We consider quantum states on $\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$, where \mathcal{H}_1 and \mathcal{H}_2 are respectively n and m ($n \leq m$) dimensional complex Hilbert spaces. We present the explicit manifold constituted by the states with certain Schmidt ranks or with given Schmidt coefficients, and calculate the dimensions of the related manifolds.

For the convenience, in the following in stead of $|x\rangle$, we simply denote x as a vector in \mathcal{H} and denote $D^1(\mathcal{H})$ as the set of all $x \in \mathcal{H}$. For any $x \in \mathcal{H}$, x can be written as the summation of tensor products,

$$x = x_1 \otimes y_1 + x_2 \otimes y_2 + \dots + x_k \otimes y_k, \quad k \in \mathbb{N}, \tag{1}$$

where $x_i \in \mathcal{H}_1, y_i \in \mathcal{H}_2$. We call the expression (??) linearly independent if x_1, x_2, \dots, x_k ; y_1, y_2, \dots, y_k are linearly independent vectors respectively. We say the length of x is k if (??) is a linearly independent expression. In fact one can easily prove that the length is just the Schmidt rank and the Schmidt decomposition is a special expression of a linearly independent one. Therefore the length of x in all linearly independent expressions is the same and the terms of tensor products contained in the linearly independent expression of x are the least in all other possible expressions of x.

[Lemma] If $x \in \mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$ has the following two linearly independent expressions

$$x = x_1 \otimes y_1 + x_2 \otimes y_2 + \dots + x_k \otimes y_k, \ x = w_1 \otimes z_1 + w_2 \otimes z_2 + \dots + w_k \otimes z_k,$$
 (2)

then there exists a non-degenerate $k \times k$ matrix C such that

$$(z_1, \dots, z_k) = (y_1, \dots, y_k)C, \quad (w_1, \dots, w_k) = (x_1, \dots, x_k)(C^t)^{-1}.$$
 (3)

[Proof] Expanding x_1, x_2, \dots, x_k to be the basis $x_1, \dots, x_k, x_{k+1}, \dots, x_n$ in \mathcal{H}_1 and y_1, y_2, \dots, y_k to be the basis $y_1, \dots, y_k, y_{k+1}, \dots, y_m$ in \mathcal{H}_2 , we have

$$w_j = \sum_{i=1}^n a_{ji} x_i , z_j = \sum_{i=1}^m b_{ji} y_i$$

for some $a_{ji}, b_{ji} \in \mathbb{C}$. Then from (??) we have

$$\sum_{i=1}^{n} \sum_{s=1}^{m} (\sum_{j=1}^{k} a_{ji} b_{js}) x_i \otimes y_s = \sum_{j=1}^{k} x_j \otimes y_j.$$
 (4)

Denote A (resp. B) the matrix with entries a_{ij} (resp. b_{ij}). As $\{x_i \otimes y_s : j = 1, 2, \dots, n; s = 1, 2, \dots, m\}$ is a basis of $\mathcal{H}_1 \bigotimes \mathcal{H}_2$, from (??) we have

$$A^t B = \begin{pmatrix} E_k & 0 \\ 0 & 0 \end{pmatrix}_{n \times m},\tag{5}$$

where E_k is the identity matrix of order k. If we rewrite A and B as block matrices $A = (A_{kk} \ A_{k,n-k}), B = (B_{kk} \ B_{k,m-k}),$ then (??) gives rise to $A_{kk}^t B_{kk} = E_k, B_{k,m-k} = 0,$

 $A_{k,n-k} = 0$. Namely, $A = (A_{kk} \quad 0), B = (B_{kk} \quad 0)$. Set $C = A_{kk}^{-1}$, we obtain $B_{kk} = (A_{kk}^t)^{-1}$ and $(w_1, \dots, w_k) = (x_1, \dots, x_k)(C^t)^{-1}, (z_1, \dots, z_k) = (y_1, \dots, y_k)B_{kk}^t = (y_1, \dots, y_k)C$.

[Theorem 1] Let $D_k^1(\mathcal{H})$, a submanifold of $D^1(\mathcal{H})$, be the set of all normalized pure states with length k, $D_k^1(\mathcal{H}) = \{x \in \mathcal{H}, \text{ the length of } x \text{ is } k, ||x||^2 = 1\}$. We have

$$D_k^1(\mathcal{H}) \simeq G(n,k) \times (CP^{k^2-1} \setminus \overline{M}) \times G(m,k),$$

where \overline{M} is a hypersurface of CP^{k^2-1} , G(n,k) is the Grassmannian manifold.

[Proof] We first prove that there is a one-to-one correspondence between $D_k^1(\mathcal{H})$ and $G(n,k)\times (CP^{k^2-1}\backslash \overline{M})\times G(m,k)$.

For $x \in D_k^1(\mathcal{H})$, suppose $x = x_1 \otimes y_1 + x_2 \otimes y_2 + \cdots + x_k \otimes y_k$ is a linearly independent expression of x. Because y_1, \dots, y_k are linearly independent, y_1, \dots, y_k span a k-dimensional subspace D_k of \mathcal{H}_2 . We fix an orthonormal basis y_1^0, \dots, y_k^0 in D_k and assume $(y_1^0 \cdots y_k^0) = (y_1 \cdots y_k)A$, where A is a non-degenerate complex $k \times k$ matrix. If we keep x unchanged, from Lemma x_1, \dots, x_k are transformed correspondingly to x_1', \dots, x_k' , $(x_1' \cdots x_k') = (x_1 \cdots x_k)(A^t)^{-1}$.

A k-dimensional subspace of \mathcal{H}_2 just corresponds to a point in a Grassmannian manifold G(m,k). As x_1', \dots, x_k' in the expression $x = x_1' \otimes y_1^0 + \dots + x_k' \otimes y_k^0$ are linearly independent, they span a k-dimensional subspace C_k of \mathcal{H}_1 . If we fix an orthonormal basis x_1^0, \dots, x_k^0 in C_k , then there exists a unique non-degenerate $k \times k$ matrix G such that $(x_1, \dots, x_k) = (x_1^0, \dots, x_k^0)G$. A k-dimensional subspace of \mathcal{H}_1 just corresponds to a point in a Grassmannian manifold G(n,k). Suppose $(x_1', \dots, x_k') = (x_1^0, \dots, x_k^0)B$, where B is a $k \times k$ complex matrix with entries b_{ij} satisfying $\sum_{i,j=1}^k |b_{ij}|^2 = 1$, $det(B) \neq 0$. Then all $B = (b_{ij})_{i,j=1}^k$ constitute a set D which can be viewed as a subset of the identity ball S^{k^2-1} in \mathbb{C}^{k^2} , where

$$S^{k^2-1} = \{(b_{11}, b_{21}, \cdots, b_{k1}, b_{12}, \cdots, b_{k2}, \cdots, b_{kk}) : \sum_{i,j=1}^{k} |b_{ij}|^2 = 1, b_{ij} \in \mathbb{C}\}.$$

Moreover, D is an open subset in S^{k^2-1} .

In summary, to determine x'_1, \dots, x'_k , we need to determine the k-dimensional subspace C_k which is spanned by x'_1, \dots, x'_k and the nondegenerate $k \times k$ matrix B associated with x'_1, \dots, x'_k , i.e. a point of Grassmannian manifold G(n, k) and a point of D are determined.

We define $A \sim B$ iff there exists $\theta \in \mathbb{R}$ such that $A = e^{i\theta}B$ and denote the equivalence class containing A by [A], then

$$S^{k^2-1}/\sim = CP^{k^2-1}$$
.

Define

$$\pi: S^{k^2-1} \longrightarrow CP^{k^2-1}$$

$$A \longrightarrow [A].$$

Then π is an open map. Suppose the image of D under π is \overline{D} which is an open subset of CP^{k^2-1} , so it is an open submanifold. Suppose $\overline{M} = CP^{k^2-1} \setminus \overline{D}$, i.e. \overline{M} is the image of the set under the map π which consists of the points contained in S^{k^2-1} satisfying det(B) = 0 and \overline{M} is a hypersurface of CP^{k^2-1} . So we have

$$\overline{D} = CP^{k^2 - 1} \backslash \overline{M}.$$

As $(e^{i\theta}x'_1, \dots, e^{i\theta}x'_k) = e^{i\theta}(x_1^0, \dots, x_k^0)B = (x_1^0, \dots, x_k^0)(e^{i\theta}B)$, the action of $e^{i\theta}$ on x can be viewed as on matrix B associated with x'_1, \dots, x'_k . So the equivalence class [x] containing x corresponds to the equivalence class [B] which contains B, i.e. [x] corresponds a point in \overline{D} . Hence, a pure state x corresponds to a unique point p in $G(n,k) \times (CP^{k^2-1} \setminus \overline{M}) \times G(m,k)$, where the coordinates of p are determined uniquely by the k-dimensional subspace D_k in \mathcal{H}_2 spanned by y_1, \dots, y_k , the k-dimensional subspace C_k in \mathcal{H}_1 spanned by x_1, \dots, x_k and [B]. We denote this kind of correspondence as F. One can easily prove that F is surjective and injective. So we get a one-to-one correspondence between $D_k^1(\mathcal{H})$ and $G(n,k) \times (CP^{k^2-1} \setminus \overline{M}) \times G(m,k)$. Moreover from the above proof we know that F is smooth.

We now imbed $G(n,k) \times (CP^{k^2-1} \setminus \overline{M}) \times G(m,k)$ to CP^{mn-1} according to F. For arbitrary $p \in G(n,k) \times (CP^{k^2-1} \setminus \overline{M}) \times G(m,k)$, the coordinates of p have the form,

$$(x_{1,k+1},\cdots,x_{2n},\cdots,x_{kn},a_{12},\cdots,a_{kk},y_{1,k+1},\cdots,y_{km}).$$

Let us write the coordinates $(x_{1,k+1}, x_{1,k+2}, \cdots, x_{kn})$ in G(n,k) in the matrix form

$$X = \begin{pmatrix} 1 & 0 & \cdots & 0 & x_{1,k+1} & \cdots & x_{1n} \\ 0 & 1 & \cdots & 0 & x_{2,k+1} & \cdots & x_{2n} \\ \vdots & \vdots & \ddots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & 1 & x_{k,k+1} & \cdots & x_{kn} \end{pmatrix},$$

and the coordinates $(a_{12}, \dots, a_{1k}, \dots, a_{kk})$ in CP^{k^2-1} in the form

$$A = \begin{pmatrix} 1 & a_{12} & \cdots & a_{1k} \\ a_{21} & a_{22} & \cdots & a_{2k} \\ \vdots & \vdots & \ddots & \vdots \\ a_{k1} & a_{k2} & \cdots & a_{kk} \end{pmatrix},$$

Set $X^tA = B$. Then $B = (b_{ij})$ is an $n \times k$ matrix. Let e_1, \dots, e_n (resp. d_1, \dots, d_m) be an orthonormal basis in \mathcal{H}_1 (resp. \mathcal{H}_2). Take

$$x_1 = \sum_{j=1}^n b_{j1}e_j, \ x_2 = \sum_{j=1}^n b_{j2}e_j, \ \cdots, \ x_k = \sum_{j=1}^n b_{jk}e_j,$$

and

$$y_1 = \sum_{j=1}^m y_{1j}d_j, \ y_2 = \sum_{j=1}^m y_{2j}d_j, \ \cdots, \ y_k = \sum_{j=1}^m y_{kj}d_j,$$

where y_{ij} are the entries of the matrix Y,

$$Y = \begin{pmatrix} 1 & 0 & \cdots & 0 & y_{1,k+1} & \cdots & y_{1m} \\ 0 & 1 & \cdots & 0 & y_{2,k+1} & \cdots & y_{2m} \\ \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & \cdots & 1 & y_{k,k+1} & \cdots & y_{km} \end{pmatrix},$$

then both x_1, \dots, x_k and y_1, \dots, y_k are linearly independent respectively.

Let

$$x = x_1 \otimes y_1 + \dots + x_k \otimes y_k = \sum_{j=1}^n \sum_{s=1}^m (\sum_{l=1}^k b_{jl} y_{ls}) e_j \otimes d_s.$$

Then $x \in D_k^1(\mathcal{H})$ is the image of p under F. Since $\mathcal{H}_1 \simeq \mathbb{C}^n$, $\mathcal{H}_2 \simeq \mathbb{C}^m$ and $D^1(\mathcal{H}) \simeq CP^{nm-1}$, we can define the imbedding:

$$f: G(n,k) \times (CP^{k^2-1} \backslash \overline{M}) \times G(m,k) \longrightarrow CP^{mn-1}$$
 $p \longrightarrow q$

where q=f(p)=x. The homogeneous coordinates of q are given by $q=(d_{11},d_{12},\cdots,d_{1m},d_{21},\cdots,d_{2m},\cdots,d_{nm})$, where $d_{js}=\sum_{l=1}^k b_{jl}y_{ls}=\sum_{l=1}^k \sum_{t=1}^k x_{tj}a_{tl}y_{ls}$ $(j=1,\cdots,n,s=1,\cdots,m)$. Then the coordinate components of q are polynomial of the coordinate components of p. Hence, the imbedding f is non-degenerate holomorphic mapping. Moreover, we have

$$f(G(n,k) \times (CP^{k^2-1} \setminus \overline{M}) \times G(m,k)) = D_k^1(\mathcal{H}).$$

Therefore $D^1_k(\mathcal{H})$ is a complex submanifold of $\mathbb{C}P^{mn-1}$ (i.e. $D^1(\mathcal{H})$), and

$$D_k^1(\mathcal{H}) \simeq G(n,k) \times (CP^{k^2-1} \setminus \overline{M}) \times G(m,k).$$

Theorem 2 The subset $D_k^1(\mu_1, \dots, \mu_k)$ of $D_k^1(\mathcal{H})$ of pure states with the Schmidt coefficients $\mu_1 \geqslant \mu_2 \geqslant \dots \geqslant \mu_k$ is a submanifold of real dimension 2k(m+n-k)-k-1, which is diffeomorphically equivalent to a manifold

$$(CP^{n-1} \times CP^{m-1}) \times \cdots \times (CP^{n-k} \times CP^{m-k}) \times T^{k-1},$$

where T^{k-1} is a torus of real dimension k-1.

[proof] For any pure state [e] of $D_k^1(\mathcal{H})$, the unit vector e has the following Schmidt representation $e = \mu_1 a_1 \otimes b_1 + \cdots + \mu_k a_k \otimes b_k$, where a_i^*s and b_i^*s are orthonormal vectors in \mathcal{H}_1 and \mathcal{H}_2 respectively, and μ_i^*s are Schmidt coefficients of e, we assume that $\mu_1 \geqslant \mu_2 \geqslant \cdots \geqslant \mu_k$. Consider the element \tilde{e} which has the same Schmidt coefficients as e, $\tilde{e} = \mu_1 \tilde{a_1} \otimes \tilde{b_1} + \cdots + \mu_k \tilde{a_k} \otimes \tilde{b_k}$, and $[a_i] = [\tilde{a_i}] \in D^1(\mathcal{H}_1)$, $[b_i] = [\tilde{b_i}] \in D^1(\mathcal{H}_2)$, $i = 1, \dots, k$. Hence \tilde{e} must have the form $\tilde{e} = \mu_1 e^{i\theta_1} a_1 \otimes b_1 + \cdots + \mu_k e^{i\theta_k} a_k \otimes b_k$, and $[\tilde{e}]$ constitute a set

$$\{([a_1], [b_1], \cdots, [a_k], [b_k], e^{i\beta_1}, \cdots, e^{i\beta_{k-1}}) \mid \beta_1, \cdots, \beta_{k-1} \in \mathbb{R}\} \simeq \mathbb{R}^{k-1}.$$

Then all the pure states with the same Schmidt coefficients $\mu_1 \geqslant \mu_2 \geqslant \cdots \geqslant \mu_k$ constitute a set which is equivalent to a manifold $(CP^{n-1} \times CP^{m-1}) \times \cdots \times (CP^{n-k} \times CP^{m-k}) \times T^{k-1}$, which is of real dimension 2k(m+n-k)-k-1.

As a simple example, let us first take $dim(\mathcal{H}_1) = dim(\mathcal{H}_2) = 3$, k = 1. For arbitrary $x \in D^1(\mathcal{H}_1)$, $y \in D^1(\mathcal{H}_2)$, by the Segre imbedding we have $Seg(x,y) = |x \otimes y\rangle\langle x \otimes y|$. As $w = x \otimes y \in D_1^1(\mathcal{H})$, one gets

$$Seg(D^1(\mathcal{H}_1) \times D^1(\mathcal{H}_2)) \subset D^1_1(\mathcal{H}).$$

And for arbitrary $w \in D_1^1(\mathcal{H})$, there exist $x \in \mathcal{H}_1$, $y \in \mathcal{H}_2$ such that $w = x \otimes y = Seg(x, y)$. The Segre imbedding $Seg : D^1(\mathcal{H}_1) \times D^1(\mathcal{H}_2) \to D_1^1(\mathcal{H})$ is a surjective map to $D_1^1(\mathcal{H})$. Hence, we have $Seg(D^1(\mathcal{H}_1) \times D^1(\mathcal{H}_2)) = D_1^1(\mathcal{H})$. Therefore $D_1^1(\mathcal{H}) \cong D^1(\mathcal{H}_1) \times D^1(\mathcal{H}_2) \cong CP^2 \times CP^2$. From Theorem 1, in this case $k^2 - 1 = 0$. We get $D_1^1(\mathcal{H}) \cong CP^2 \times CP^2$.

As a more complicated case, we consider $dim(\mathcal{H}_1) = 3$, $dim(\mathcal{H}_2) = 4$, and k = 2. For arbitrary $p \in G(3,2) \times (CP^3 \setminus \overline{M}) \times G(4,2)$ with coordinate $p = (x_{13}, x_{23}, a_{12}, a_{21}, a_{22}, y_{13}, y_{14}, y_{23}, y_{24})$, set

$$X = \begin{pmatrix} 1 & 0 & x_{13} \\ 0 & 1 & x_{23} \end{pmatrix}, \quad A = \begin{pmatrix} 1 & a_{12} \\ a_{21} & a_{22} \end{pmatrix}, \quad Y = \begin{pmatrix} 1 & 0 & y_{13} & y_{14} \\ 0 & 1 & y_{23} & y_{24} \end{pmatrix}.$$

Then

$$X^{t}A = \begin{pmatrix} 1 & a_{12} \\ a_{21} & a_{22} \\ x_{13} + a_{21}x_{23} & a_{12}x_{13} + x_{23}a_{22} \end{pmatrix}.$$

We take $x_1 = e_1 + a_{21}e_2 + (x_{13} + a_{21}x_{23})e_3$, $x_2 = a_{12}e_1 + a_{22}e_2 + (a_{12}x_{13} + x_{23}a_{22})e_3$, $y_1 = d_1 + y_{13}d_3 + y_{14}d_4$, $y_2 = d_2 + y_{23}d_3 + y_{24}d_4$. Let $x = x_1 \otimes y_1 + x_2 \otimes y_2$. Since $x_1, x_2; y_1, y_2$ are linearly independent respectively, we have $x \in D_2^1(\mathcal{H})$.

For arbitrary $x \in D_2^1(\mathcal{H})$, suppose $x = x_1 \otimes y_1 + x_2 \otimes y_2$, then y_1, y_2 span a unique 2-dimensional subspace D_2 of \mathcal{H}_2 . We fix an orthonormal basis y_1^0, y_2^0 in D_2 and suppose $(y_1^0, y_2^0) = (y_1, y_2)A$, where A is a non-degenerate complex 2×2 matrix. At the same time, suppose that x_1, x_2 are transformed correspondingly to $x_1', x_2', (x_1', x_2') = (x_1, x_2)(A^t)^{-1}$. Then $x = x_1' \otimes y_1^0 + x_2' \otimes y_2^0$, and x_1', x_2' generate a unique 2-dimensional subspace C_2 of \mathcal{H}_1 . We fix an orthonormal basis x_1^0, x_2^0 in C_2 and assume $(x_1', x_2') = (x_1^0, x_2^0)B$. Then (x_1', x_2') are determined uniquely by C_2 and B. Moreover, $[x_1', x_2']$ correspond to [B], and $[B] \in \mathbb{C}P^3 \setminus \overline{M}$, where $\overline{M} = \{[A]: A$ are complex 2×2 matrices, $det(A) = 0\}$. C_2 is associated to a point of Grassmannian manifold G(3, 2) and D_2 is associated to a point of Grassmannian manifold G(3, 2), i.e. x is associated to a point of $G(3, 2) \times (\mathbb{C}P^3 \setminus \overline{M}) \times G(4, 2)$.

Furthermore, for arbitrary point in $G(3,2) \times (CP^3 \setminus \overline{M}) \times G(4,2)$, we can find correspondingly a unique point in $D_2^1(\mathcal{H})$, and vice versa. In this case, the imbedding is

$$f: \ G(3,2)\times (CP^3\backslash \overline{M})\times G(4,2) \ \longrightarrow \ CP^{11}$$

$$p \ \longrightarrow \ q$$

where q = f(p) = x and the homogeneous coordinates of q are assumed to be $q = (d_{11}, d_{12}, d_{13}, d_{14}, d_{21}, d_{22}, d_{23}, d_{24}, d_{31}, d_{32}, d_{33}, d_{34})$, where $d_{11} = 1, d_{12} = a_{12}, d_{13} = a_{12}y_{23} + y_{13}, d_{14} = a_{12}y_{24} + y_{14}, d_{21} = a_{21}, d_{22} = a_{22}, d_{23} = a_{21}y_{13} + a_{22}y_{23}, d_{24} = a_{21}y_{14} + a_{22}y_{24}, d_{31} = x_{13} + a_{21}x_{23}, d_{32} = a_{12}x_{13} + a_{22}x_{23}, d_{33} = y_{23}(a_{12}x_{13} + a_{22}x_{23}) + y_{13}(x_{13} + a_{21}x_{23}), d_{34} = y_{24}(a_{12}x_{13} + a_{22}x_{23}) + y_{14}(x_{13} + a_{21}x_{23}).$

The first example tests the theorem from the Segre imbedding point of view. In this case the second factor of the product manifold generates a point. The second one is a lower dimension case according to the Theorem 1.

We have investigated the complex manifold structure for bipartite pure states and the Kähler metric of $D_k^1(\mathcal{H})$, by presenting explicitly the manifolds with respect to all pure states

of arbitrarily given Schmidt ranks or Schmidt coefficients and calculating the dimensions of the corresponding manifolds. In fact, we also can express the Kähler metric of $D_k^1(\mathcal{H})$ by local coordinates, but the expressions are very complicated and it is difficult to compute the geometrical objects such as holomorphic curvature, scalar curvature. It would be also nice to describe the entanglement of quantum states according to some functions of metric or geometrical objects. The results in this paper can be used to study the differential geometry of bipartite mixed states.

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^[] M.A. Nielsen and I.L. Chuang, Quantum Computation and Quantum Information, Cambridge University Press, Cambridge, 2000.

^[] C.H. Bennett, D.P. DiVincenzo, J.A. Smolin, and W.K. Wootters, Phys. Rev. A 54, 3824 (1996).

M. Horodecki, Quant. Inf. Comp. 1, 3 (2001).

D. Bruß, J. Math. Phys. 43, 4237 (2002).

W.K. Wootters, Phys. Rev. Lett. 80, 2245(1998).

B.M. Terhal and K.G.H. Vollbrecht, Phys. Rev. Lett. 85, 2625 (2000).

S.M. Fei and X.Q. Li-Jost, Rep. Math. Phys. 53, 195 (2004).

K. Chen, S. Albeverio, and S.M. Fei, Phys. Rev. Lett. 95, 040504 (2005).

K. Chen, S. Albeverio, and S.M. Fei, Phys. Rev. Lett. 95, 210501 (2005). S.M. Fei, Z.X. Wang and H. Zhao, Phys. Lett. A 329(2004)414-419.

^[] M. Lewenstein, D. Bruß, J. I. Cirac, B. Kraus, M. Kuś, J. Samsonowicz, A. Sanpera, and R. Tarrach, J. Mod. Phys. 47, 2481 (2000).

R. Werner, Phys. Rev. A40, 4277 (1989).

A. Peres Phys. Rev. Lett. 77, 1413 (1996).

M. Horodecki, P. Horodecki, and R. Horodecki, Phys. Lett. A 223, 8 (1996).

W. Dür, G. Vidal and J.I. Cirac, Phys. Rev. A 62, 062314(2000).

- S. Karnas and M. Lewenstein, Phys. Rev. A 64, 042313 (2001).
- S. Albeverio, S.M. Fei and D. Goswami, Phys. Lett. A 286 (2001)91-96.
- S.M. Fei, X.H. Gao, X.H. Wang, Z.X. Wang and K. Wu, Phys. Lett. A 300 (2002)559-566;Phys. Rev. A 68 (2003) 022315.
- S. Albeverio, K. Chen and S.M. Fei, Phys. Rev. A 68(2003)062313.
- [] E.M. Rains, IEEE Transactions on Information Theory 46 54-59(2000).
 - M. Grassl, M. Rötteler and T. Beth, *Phys. Rev. A* 58, 1833(1998).
 - Y. Makhlin, Quant. Info. Proc. 1, 243-252 (2002).
 - N. Linden, S. Popescu and A. Sudbery, Phys. Rev. Lett. 83, 243 (1999).
 - S. Albeverio, S.M. Fei, P. Parashar and W.L. Yang, Phys. Rev. A 68 (Rapid Comm.) (2003) 010303.
 - S. Albeverio, L. Cattaneo, S.M. Fei and X.H. Wang, Int. J. Quant. Inform. **3**(2005)603-609; Rep. Math. Phys. **56** (2005)341-350.
 - S.M. Fei and N.H. Jing, Phys. Lett. A 342(2005)77-81.
 - C. H. Bennett, S. Popescu, D. Rohrlich, J. A. Smolin, A. V. Thapliyal, Phys. Rev. A 63 (2001) 012307.
 - W. Dür, G. Vidal, J. I. Cirac, Phys. Rev. A 62, 062314(2000).
 - F. Verstraete, J.Dehaene, B.De Moor and H. Verschelde Phys. Rev. A. 65, 052112 (2002).
- [] Janusz Grabowski, Giuseppe Marmo, Maret Kus, Geometry of quantum systems:density states and entanglement, J. phys. A 38(2005)10217-10244.
- [] V. I. Man'ko, G. Marmo, E. C. G. Sudarshan, F. Zaccaria, Differential geometry of density states, Rept. Math. Phys. 55(2005)405-422.
- [] J. Grabowski, M. Kuś and G. Marmo, Open Sys. and Information Dyn. 13(2006)343-362.